

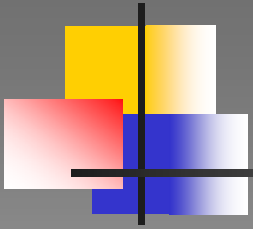
Solving the η -problem in Hybrid Inflation with Heisenberg Symmetry and Stabilized Modulas

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Dark Matter at the Crossroads, DESY, Hamburg



Based on - [arXiv:0808.2425 \[hep-ph\]](https://arxiv.org/abs/0808.2425)
with S. Antusch, M. Bastero-Gil, S.F. King, P. Kostka



Summary

- A class of SUSY hybrid inflationary model where η -problem is resolved by Heisenberg symmetry
- Associated modulus gets stabilized by large vacuum energy during inflation
- At tree level, inflaton direction is flat - protected by symmetry
- Radiative corrections lift the flatness such that $n_s < 1$ and consistent with data

Motivation

- Cosmology is now a precision science



Kolb and Turner

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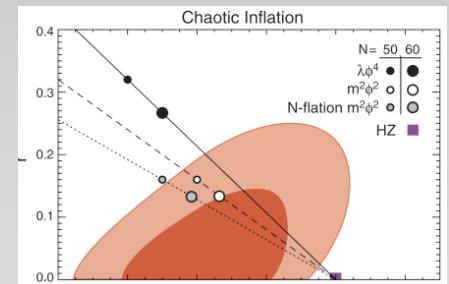
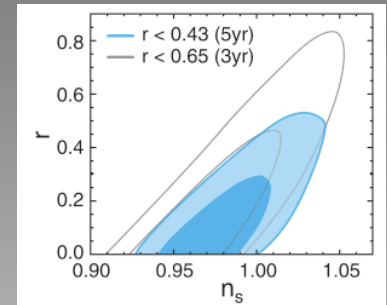
Dodelson

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Mukhanov

- Inflation is a widely accepted paradigm for early Universe physics, but seeking for a consistent particle physics model
- Data is not precise enough and dynamics is not well understood
- Slow-Roll conditions are difficult to reconcile with known interactions of particle physics
- We present one possible toy model





η -problem

- Consider a flat inflation potential $V(\phi)$

no reason to forbid the operator $\bar{\phi}\phi \frac{V(\phi)}{M_P^2}$

$$\eta = \left(\frac{V''}{V} \right) \sim 1$$

- In SUSY models effects of low energy SUGRA important
flat direction ($\eta \ll 1$) \rightarrow extremely curved ($\eta \sim 1, m_\phi \sim H$)

$$V(\phi) = e^{K(\phi)/M_P^2} V_0 \sim V_0 \left(1 + \frac{\phi^2}{M_P^2} + \dots \right) \sim V_0 + 3H^2 \phi^2$$

Copeland, Liddle, Lyth, Stewart, Wands and Dine, Randall, Thomas



Moduli Problem

- ❑ Several extra scalar fields - needs to make them heavy (Dine)
- ❑ Scalar fields must find a stable minimum (Brustein, Steinhardt)
- ❑ Initial assumption: Moduli stabilization has no effect on inflationary dynamics - **WRONG!**
(Brax, van de Bruck, Davis, Davis)
- ❑ Problem: Requirements for moduli and inflaton are exactly opposite!
- ❑ Avenues for solutions:
 - a) specific choice of Kahler potential (Murayama et.al hep-ph/9311326)
 - b) symmetry requirement of Kahler potential
e.g shift symmetry or Heisenberg symmetry (our approach)

(Kawasaki et.al hep-ph/0004243, Brax and Martin, Davis and Postma)



General Framework

- Hybrid Inflation - Inflaton field N and Waterfall field H

$$W = \kappa S (g_1(H, N) - M^2) + g_2(H, N)$$

$$K = (|S|^2 + |H|^2 + \kappa_S |S|^4 + \kappa_{SH} |S|^2 |H|^2 + \dots) + g_3(\rho) |S|^2 + f(\rho)$$

- Modulus is defined by $\rho = T + T^* - |N|^2$ and invariant under Heisenberg symmetry

- Main feature: $W = 0, W_N = W_H = W_T = 0, W_S \neq 0$ but $H = S = 0$

- For comparison: the standard SUSY hybrid inflation is implemented by $W = \kappa N (\bar{H}H - M^2)$



General Framework

- g_1 has to be chosen such that $|F_S|^2 = |g_1(H = 0, N) - M^2|^2$
typically it would depend on H
- g_2 leads to positive N -dependent mass² for H via $|F_H|^2$ during inflation - $g_2 \sim N^m H^2$ with $m \geq 1$
- g_3 together with $f(\rho)$ shapes the potential for ρ
 $g_3(\rho)|S|^2$ term in the Kahler potential induce a contribution to the potential of the order of vacuum energy and effectively stabilize ρ



Explicit Example

- For rest of the discussion: $W = \kappa S (H^2 - M^2) + \frac{\lambda}{M_*} N^2 H^2$

$$K(H, S, N, T) \equiv |H|^2 + (1 + \kappa_S |S|^2 + \kappa_\rho \rho) |S|^2 + f(\rho)$$

- First term is similar to the SUSY hybrid
but $S = 0$ both during and after inflation
 $H = 0$ during inflation and $H \sim M$ in the end

$$\mathcal{L}_{\text{kin}} = \frac{f''(\rho)}{4} n^2 (\partial_\mu n)^2 - \frac{f'(\rho)}{2} (\partial_\mu n)^2 - \frac{f''(\rho)}{2\sqrt{2}} n \partial_\mu n \partial^\mu t + \frac{f''(\rho)}{8} (\partial_\mu t)^2$$

$$\mathcal{L}_{\text{kin}} = \frac{f''(\rho)}{4} (\partial_\mu \rho)^2 - \frac{f'(\rho)}{2} (\partial_\mu n)^2$$

$$V_F = e^K \left[K^{i\bar{j}} \mathcal{D}_i W \mathcal{D}_{\bar{j}} W^* - 3|W|^2 \right]$$

Analysis

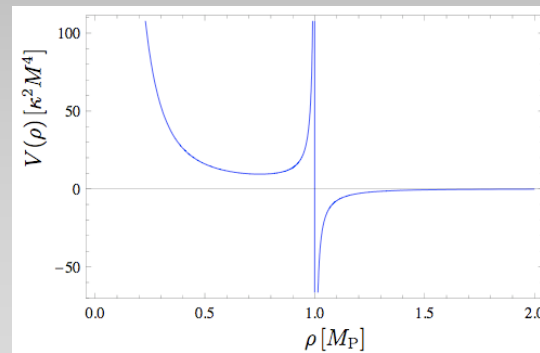
$$V_F = e^K \left[K^{i\bar{j}} \mathcal{D}_i W \mathcal{D}_{\bar{j}} W^* - 3|W|^2 \right]$$

$$V_{\text{tree}} = V_F = e^{f(\rho)} K_{SS^*}^{-1} \left| \frac{\partial W}{\partial S} \right|^2 = \kappa^2 M^4 \cdot \frac{e^{f(\rho)}}{(1 + \kappa_\rho \rho)} \quad \text{with } S=H=0$$

- Tree level potential is exactly flat in n-direction
Heisenberg symmetry protects n from obtaining large mass corrections from SUGRA expansion

- with $f(\rho) = -3 \ln(\rho)$

$$\rho_{\text{min}} = -\frac{3}{4 \kappa_\rho}$$





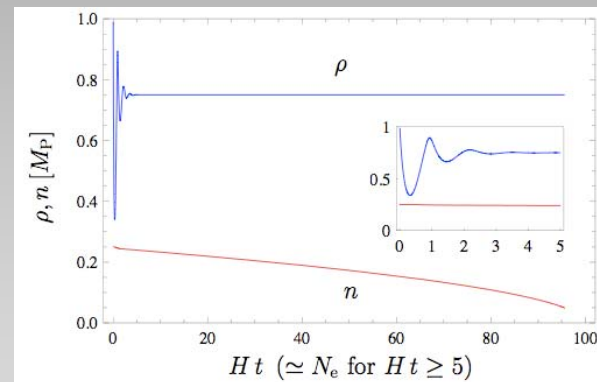
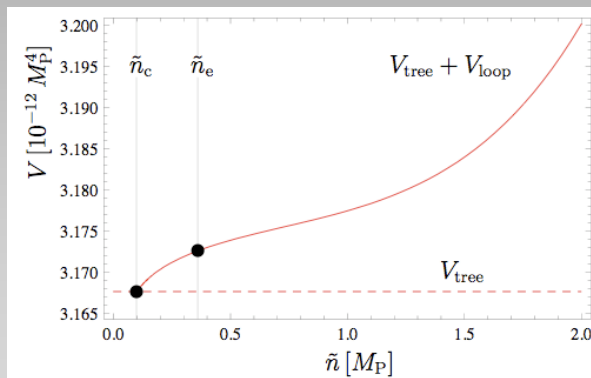
Analysis

- Without assuming $S=H=0$, we can show that in this stable patch all the masses ρ and S mass is larger than the Hubble's constant.
- Critical value of the waterfall field: $\tilde{n}_c^2 = 8 \frac{\kappa}{\lambda} (MM_*) \sqrt{2 - M^2}$
- Imaginary part of the fields redshifts away quickly and decouple from the real part

One-Loop Effective Potential

- N is exactly flat at the classical level
- Quantum corrections are generated by Heisenberg symmetry breaking W and broken SUSY during inflation
- Only h contributes to the n -dependent mass

$$V_{\text{eff}}(n, \rho) = V_{\text{tree}}(\rho) + V_{\text{loop}}(n, \rho)$$



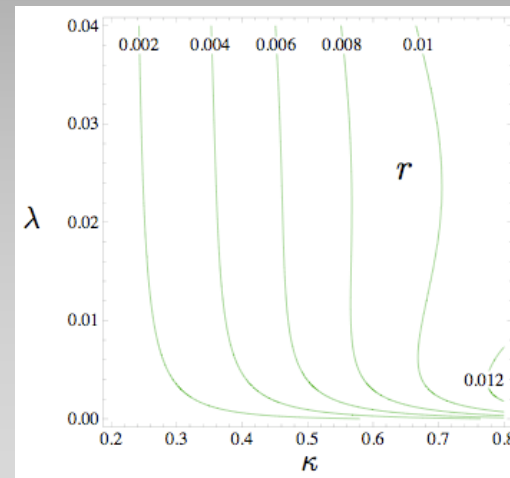
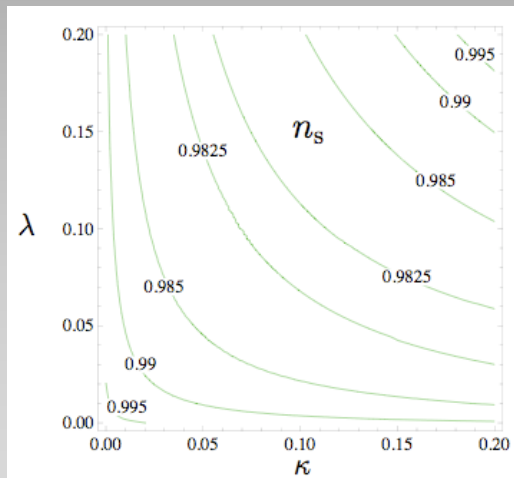
Predictions

- Predictions are very similar to the standard hybrid models

For example, $(\kappa, \lambda) = (0.05, 0.2)$

$$n_s \simeq 0.982, \quad \frac{M}{M_{\text{P}}} \simeq 3.4 \cdot 10^{-3}$$

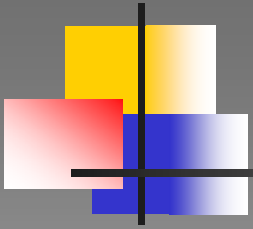
$$r \simeq 9.0 \cdot 10^{-5}, \quad \frac{dn_s}{d \ln k} \simeq -2.4 \cdot 10^{-3}$$



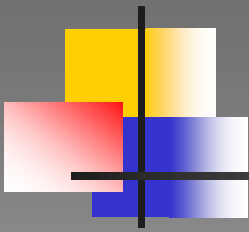


Summary and Outlook

- ❑ A class of SUSY hybrid inflationary model where η -problem is resolved by Heisenberg symmetry
- ❑ Associated modulus gets stabilized by large vacuum energy during inflation
- ❑ At tree level, inflaton direction is flat - protected by symmetry
- ❑ Radiative corrections lift the flatness such that $n_s < 1$ and consistent with data
- ❑ Right handed sneutrino is an ideal candidate for inflaton in this set up
- ❑ Embedding the model in GUT framework



Thank you



$$\mathcal{L}_{\text{SUGRA}} \supset -\frac{1}{2} m_{3/2} \left(G_{ij} + G_i G_j - G_{ij\bar{k}} G^{\bar{k}} \right) \chi_i \chi_j - \text{H.c.}$$

$$(\mathcal{M}_{\text{F}})_{ij} = e^{K/2} \left(W_{ij} + K_{ij} W + K_i W_j + K_j W_i + K_i K_j W - K^{k\bar{l}} K_{ij\bar{l}} \mathcal{D}_k W \right)$$

$$(K_{i\bar{j}}) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 + \kappa_\rho \rho + 4\kappa_S |S|^2 & -\kappa_\rho N S^* & \kappa_\rho S^* \\ 0 & -\kappa_\rho N^* S & f''(\rho) |N|^2 - f'(\rho) - \kappa_\rho |S|^2 & -f''(\rho) N^* \\ 0 & \kappa_\rho S & -f''(\rho) N & f''(\rho) \end{pmatrix}$$

$$\begin{aligned} \ddot{h} + 3H(t)\dot{h} + \frac{\partial V}{\partial h} &= 0, \\ \ddot{n} + 3H(t)\dot{n} + \frac{f''(\rho)}{f'(\rho)} \dot{\rho} \dot{n} - \frac{1}{f'(\rho)} \frac{\partial V}{\partial n} &= 0, \\ \ddot{\rho} + 3H(t)\dot{\rho} + \frac{f^{(3)}(\rho)}{2f''(\rho)} \dot{\rho}^2 + \dot{n}^2 + \frac{2}{f''(\rho)} \frac{\partial V}{\partial \rho} &= 0. \end{aligned}$$

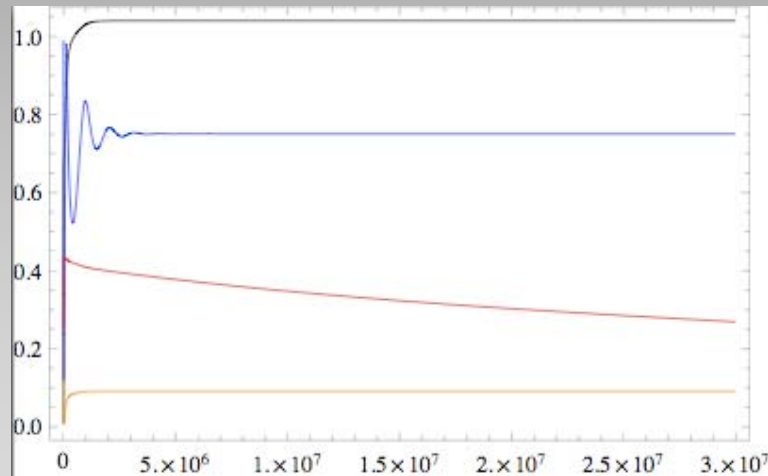
$$m_{\tilde{\rho}}^2 = -\frac{16384}{81} \kappa_{\rho}^5 \kappa^2 M^4,$$

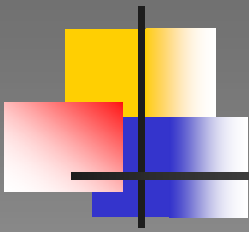
$$m_{\tilde{s}}^2 = \frac{4096}{27} \kappa_{\rho}^3 \kappa_S \kappa^2 M^4,$$

$$m_{\tilde{h}}^2 = -\frac{64}{27} \kappa_{\rho}^3 \left[\frac{\lambda^2}{16 M_*^2} \tilde{n}^4 + 8 \kappa^2 M^2 \left(\frac{M^2}{2} - 1 \right) \right]$$

$$m_{\tilde{n}}^2 = 0.$$

$$H^2 = \frac{1}{3} V(\rho_{\min}) \Big|_{s=h=0} = -\frac{256}{81} \kappa_{\rho}^3 \kappa^2 M^4$$





$$V_{\text{loop}}(n) = \frac{1}{32\pi^2} \text{Str} [\mathcal{M}^2(\tilde{n}) Q^2] + \frac{1}{64\pi^2} \text{Str} \left[\mathcal{M}^4(\tilde{n}) \left(\ln \left(\frac{\mathcal{M}^2(\tilde{n})}{Q^2} \right) - \frac{3}{2} \right) \right]$$

$$x := \left(\frac{\lambda}{\kappa} \right)^2 \frac{1 + \kappa_\rho \rho}{2 (MM_*)^2} \tilde{n}^4$$

$$m_{\text{B}}^2 = 2 \frac{(\kappa M)^2}{(1 + \kappa_\rho \rho)} e^{f(\rho)} \left[x + \frac{M^2}{2} \mp 1 \right]$$

$$m_{\text{F}}^2 = 2 \frac{(\kappa M)^4}{(1 + \kappa_\rho \rho)} e^{f(\rho)} x$$

$$V_{\text{loop}}(x) = \frac{(\kappa M)^4}{64 (1 + \kappa_\rho \rho)^2 \pi^2} \left[\begin{aligned} & 4 e^{2f(\rho)} (x + M^2/2 - 1)^2 \left[\ln \left(\frac{2 \kappa^2 M^2 e^{f(\rho)}}{(1 + \kappa_\rho \rho) Q^2} \right) + \ln (x + M^2/2 - 1) - 3/2 \right] \\ & + 4 e^{2f(\rho)} (x + M^2/2 + 1)^2 \left[\ln \left(\frac{2 \kappa^2 M^2 e^{f(\rho)}}{(1 + \kappa_\rho \rho) Q^2} \right) + \ln (x + M^2/2 + 1) - 3/2 \right] \\ & - 8 e^{2f(\rho)} x^2 \left[\ln \left(\frac{2 \kappa^2 M^2 e^{f(\rho)}}{(1 + \kappa_\rho \rho) Q^2} \right) + \ln (x) - 3/2 \right] \\ & + \frac{(\kappa M Q)^2}{16 (1 + \kappa_\rho \rho) \pi^2} \left[e^{f(\rho)} M^2 \right]. \end{aligned} \right]$$